

Partial two-particle equations in the relativistic configurational representation for scattering p -states in the case of a superposition of two δ -function potentials

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Abstract

Exact solutions of two-particle equations are considered in the relativistic configurational representation for scattering p -states in the case of superposition of two δ -function potentials. Scattering amplitudes and cross sections are determined. Comparative analysis is carried out with the case of scattering s -states.

Introduction

Partial wave equations in the relativistic configurational representation for scattering p -states (orbital angular momentum $l = 1$) of system of two particles are defined as [1]

$$\psi_{(j)}(r) = m \operatorname{sh} \chi_q p_1(\chi_q, r) + \int_0^{\infty} G_{(j)}(r, r') V(r') \psi_{(j)}(r') dr', \quad (1)$$

$$p_1(\chi_k, r) = r(mr + i)^{-1} \operatorname{sh}^{-2} \chi_k \left(\operatorname{ch} \chi_k \sin(\chi_k mr) (mr)^{-1} - \operatorname{sh} \chi_k \cos(\chi_k mr) \right). \quad (2)$$

In expressions (1) – (2) $\psi_{(j)}(r)$ – is the wave function, χ_q – is the rapidity connected with the system energy by the formula $2E = 2m \operatorname{ch} \chi_q$, m – is the mass of each particle, and the index j indicates the corresponding equation ($j=1$ – the modified Kadyshevsky equation, $j=2$ – the Logunov-Tavkhelidze equation, $j=3$ – the modified the Logunov-Tavkhelidze equation, $j=4$ – the Kadyshevsky equation). The explicit form of the partial Green's functions (GF) $G_{(j)}(r, r')$ was defined earlier [2].

When $r \rightarrow \infty$ the GF $G_{(j)}(r, r')$ take the following form:

$$G_{(j)}(r, r') \Big|_{r \rightarrow \infty} \cong K_{(j)}(r') e^{i\chi_q m r}, \quad (3)$$

$$K_{(1)}(r') = K_{(3)}(r') = -ir' \left((mr' - i) \operatorname{sh} \chi_q \right)^{-1} \left(\cos(\chi_q m r') - (mr')^{-1} \operatorname{cth} \chi_q \sin(\chi_q m r') \right),$$

$$K_{(2)}(r') = K_{(4)}(r') = -ir' (mr' - i)^{-1} \left(2 \cos(\chi_q m r') (\operatorname{sh}(2\chi_q))^{-1} - \sin(\chi_q m r') (mr' \operatorname{sh}^2 \chi_q)^{-1} \right).$$

1. Solution of the equations for a superposition of two δ -function potentials

For the following superposition of two δ -function potentials:

$$V(r) = V_{01} \delta(r - a_1) + V_{02} \delta(r - a_2) \quad (4)$$

equation (1) takes the following form:

$$\psi_{(j)}(r) = m \operatorname{sh} \chi_q p_1(\chi_q, r) + G_{(j)}(r, a_1) V_{01} \psi_{(j)}(a_1) + G_{(j)}(r, a_2) V_{02} \psi_{(j)}(a_2). \quad (5)$$

In order to determine the values of $\psi_{(j)}(a_1)$ and $\psi_{(j)}(a_2)$ in (5), it is necessary to consider the wave function at the points $r = a_1$ and $r = a_2$, and then solve the resulting system of linear algebraic equations. Expressions for $\psi_{(j)}(a_1)$ and $\psi_{(j)}(a_2)$ can be then represented as

$$\psi_{(j)}(a_1) = \Delta_{1(j)}(\chi_q) / \Delta_{(j)}(\chi_q), \quad \psi_{(j)}(a_2) = \Delta_{2(j)}(\chi_q) / \Delta_{(j)}(\chi_q); \quad (6)$$

$$\Delta_{(j)}(\chi_q) = \prod_{n=1}^2 \left[1 - V_{0n} G_{1(j)}(\chi_q, a_n, a_n) \right] - V_{01} V_{02} \left(G_{1(j)}(\chi_q, a_1, a_2) \right)^2,$$

$$\Delta_{1(j)}(\chi_q) = m \operatorname{sh} \chi_q p_1(\chi_q, a_1) \left[1 - V_{02} G_{1(j)}(\chi_q, a_2, a_2) \right] + V_{02} m \operatorname{sh} \chi_q p_1(\chi_q, a_2) G_{1(j)}(\chi_q, a_1, a_2),$$

$$\Delta_{2(j)}(\chi_q) = m \operatorname{sh} \chi_q p_1(\chi_q, a_2) \left[1 - V_{01} G_{1(j)}(\chi_q, a_1, a_1) \right] + V_{01} m \operatorname{sh} \chi_q p_1(\chi_q, a_1) G_{1(j)}(\chi_q, a_2, a_1).$$

Taking into account the asymptotic expressions for GF (3), we obtain

$$\psi_{(j)}(r) \Big|_{r \rightarrow \infty} = m \operatorname{sh} \chi_q \left(p_1(\chi_q, r) \right) \Big|_{r \rightarrow \infty} + f_{1(j)}(\chi_q) (-i) e^{im\chi_q r} m \operatorname{sh} \chi_q. \quad (7)$$

In this case the relativistic scattering amplitude $f_{1(j)}(\chi_q)$ corresponding to the equation with the model potential (4) can be expressed analytically:

$$f_{1(j)}(\chi_q) = i \left(m \operatorname{sh} \chi_q \right)^{-1} \left(K_{(j)}(a_1) V_{01} \psi_{(j)}(a_1) + K_{(j)}(a_2) V_{02} \psi_{(j)}(a_2) \right). \quad (8)$$

The partial scattering cross section for the p -wave is expressed by the scattering amplitude $f_{1(j)}(\chi_q)$ as follows: $\sigma_{1(j)}(\chi_q) = 12\pi \left| f_{1(j)}(\chi_q) \right|^2$. Figure 1 shows the graphs of the dependence of the scattering cross sections on the rapidity for the following parameter values: $a_1 = 3$, $V_{01} = 1$, $a_2 = 4$, $V_{02} = 2$, $m = 1$.

2. Comparison of results for scattering s- and p-states

Earlier [3], the scattering cross sections were obtained for the equations under consideration in the case of an s-wave. Let us give the graphs of the scattering cross section for the Logunov-

Tavkhelidze equation in the case of $l = 0.1$ for $a_1 = 3$, $V_{01} = 1$, $a_2 = 4$, $V_{02} = 2$, $m = 1$ (Figure 2). Analyzing Fig. 2, we see that the local maxima of the scattering cross section for the p -state compared to the maxima for the s -state are shifted in the direction of high rapidity values. This feature is a consequence of the presence of a centrifugal barrier.

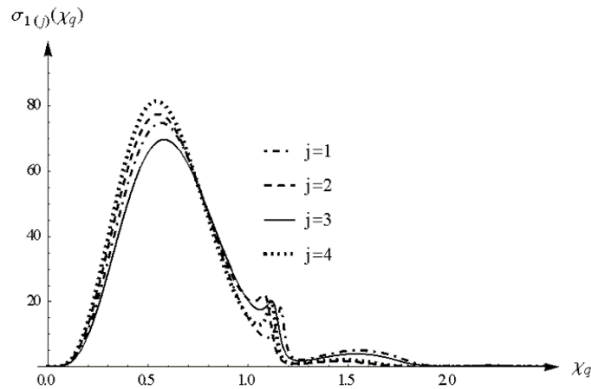


Fig. 1. Dependence of the partial scattering cross section $\sigma_{1(j)}$ on χ_q

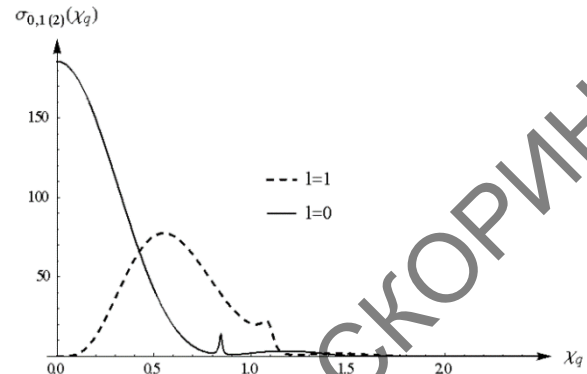


Fig. 2. Dependence of the partial scattering cross section $\sigma_{0,1(2)}$ on χ_q

Conclusion

Thus, in this paper, we consider the determination of the wave functions of partial two-particle equations in the relativistic configurational representation in case of a unit orbital angular momentum. On the basis of the obtained wave functions the partial scattering amplitudes and scattering cross sections for the superposition of two δ -function potentials are found.

References

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